

SCIENTIFIC REPORTS



OPEN

Thermoelectric Power in Bilayer Graphene Device with Ionic Liquid Gating

Yung-Yu Chien¹, Hongtao Yuan², Chang-Ran Wang¹ & Wei-Li Lee¹

Received: 24 November 2015

Accepted: 06 January 2016

Published: 08 February 2016

The quest for materials showing large thermoelectric power has long been one of the important subjects in material science and technology. Such materials have great potential for thermoelectric cooling and also high figure of merit ZT thermoelectric applications. We have fabricated bilayer graphene devices with ionic-liquid gating in order to tune its band gap via application of a perpendicular electric field on a bilayer graphene. By keeping the Fermi level at charge neutral point during the cool-down, we found that the charge puddles effect can be greatly reduced and thus largely improve the transport properties at low T in graphene-based devices using ionic liquid gating. At $(V_{ig}, V_{bg}) = (-1V, +23V)$, a band gap of about 36.6 ± 3 meV forms, and a nearly 40% enhancement of thermoelectric power at $T = 120$ K is clearly observed. Our works demonstrate the feasibility of band gap tuning in a bilayer graphene using ionic liquid gating. We also remark on the significant influence of the charge puddles effect in ionic-liquid-based devices.

Graphene-based device remains to be one of the most important subjects in both fundamental and applied researches. In fundamental sciences, many intriguing phenomena were discovered further extending the scope of our knowledge on material sciences, such as the quantum spin Hall phase in bilayer graphene^{1,2} and transverse magnetic focusing of electrons at room temperature in few layer graphene³. For practical applications, a number of experiments using graphene-based plasmonic and photonic devices draw many attentions recently⁴⁻⁶. Several optoelectronics devices have been proposed and demonstrated in graphene-based nanostructures, where fast electro-optics response and high tunability make it a wonderful platform for the application consideration. One important characteristic in a bilayer graphene (BLG) is the band gap tunability via a perpendicular electric field. A sizable band gap can be induced in a BLG by either chemical doping⁷ or dual-gated field effect⁸ to break the inversion symmetry of a BLG. It is, therefore, a challenging task to achieve an even bigger band-gap in graphene-based device, where higher on-off ratio and also thermoelectric power (TEP) can be largely enhanced⁹.

According to the Mott relation¹⁰, TEP of a material is proportional to the density of states (DOS) at Fermi level. In principle, if a giant band gap can be induced in a BLG, one expects to see a much larger DOS at the band edges leading to a huge TEP enhancement. In an earlier work, a four-fold enhancement in TEP at low temperature was demonstrated in a dual-gated BLG device using hexagonal boron nitride as a top gate dielectric material⁹. Unfortunately, the highest electric field can be applied on a BLG is only limited to about 1 V/nm mostly constrained by the breakdown voltage of the dielectric materials. One possible solution is to use high κ dielectric materials, such as HfO₂. Nevertheless, deposition of oxide material on top of few layer graphene is known to greatly reduce the carrier mobility¹¹. It is, therefore, an important subject to look for alternative approaches to fabricate top gate in graphene-based devices without degrading the graphene quality. In this report, we explored such a possibility by utilizing the ionic liquid (DEME-TFSI) gating technique in BLG devices. Ionic liquid is a molten salt at room temperature. Above its freezing point, the ions are mobile and can respond to an applied electric field. The ionic liquid side gate can serve as a top gate to introduce a large electric field on the top layer of BLG owing to the formation of the electric double layer (EDL)^{12,13}. The thickness of EDL is typically around 1 nm, which provides a strong electric field gradient on the top surface of a material. The surface carrier density can be boosted up to the order of 10^{14} per cm² that is difficult to achieve using regular oxide gating method. By controlling both the ionic liquid gate voltage and the bottom gate voltage, a sizable band-gap opening effect has been reported in a BLG device using ionic liquid gating^{14,15}. Nevertheless, the associated TEP measurements remain lacking.

¹Institute of Physics, Academia Sinica, Nankang, Taipei 11529, Taiwan. ²Geballe Laboratory for Advanced Materials, Stanford University, Stanford, California 94305, USA. Correspondence and requests for materials should be addressed to W.-L.L. (email: wlee@phys.sinica.edu.tw)

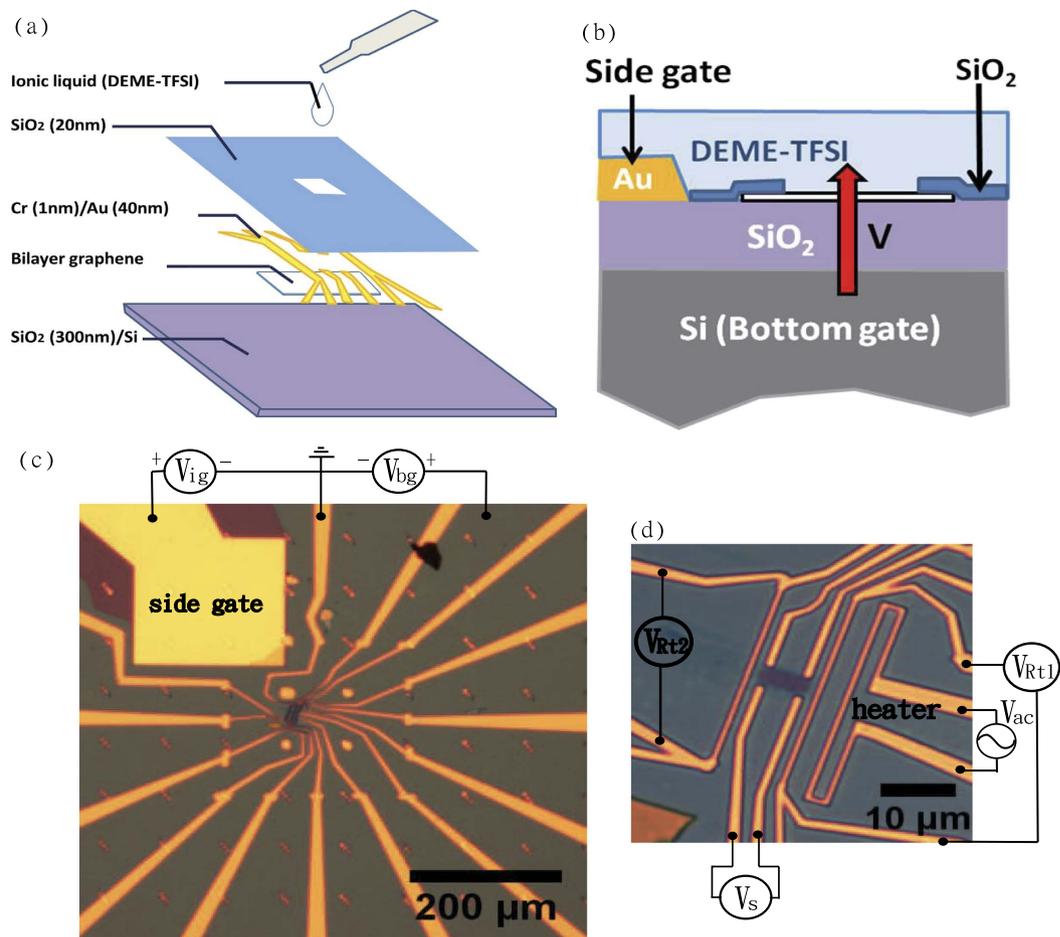


Figure 1. (a) An illustration of fabrication procedure for BLG device with ionic liquid gating. (b) A cross-sectional view of an ionic liquid gating BLG device. (c) An optical image at lower magnification. The large gold pad on the upper-left corner is the side gate for ionic liquid gating. (d) An optical image of a BLG device with ionic liquid gating. The gold contact pads are designed for both the TEP and charge transport measurements.

Results and Discussions

Figure 1(a) illustrates the fabrication procedure for a BLG device with ionic liquid gating. First, a microcrystal of BLG was mechanically exfoliated on a SiO₂ (300 nm)/Si substrate followed by the gold contact pads fabrication using e-beam lithography. A thin layer of SiO₂ (≈ 20 nm) were then evaporated on the BLG device to protect the device from the direct contact with the ionic liquid. Only a small window on the BLG was directly in contact with the ionic liquid. The cross-sectional view of the resulting device is shown in Fig. 1(b), where a side gate via the DEME-TFSI ionic liquid acts as a top gate. Both the side gate and bottom gate work simultaneously giving an electric field perpendicular to the BLG. Figure 1(d) shows an optical image of a BLG device we fabricated. The small rectangle at the center is the opening window to the BLG, where ionic liquid is directly in contact with the BLG. Other parts of the device were covered by a thin layer of SiO₂ with thickness of about 20 nm. The contact pads of Ti/Au were fabricated by electron beam lithography, which comprises a local heater and two local thermometers (V_{Rt1} , V_{Rt2}) and 4 voltage leads on a BLG for electrical signal (V_s) detection. The side gate and bottom gate are shown in Fig. 1(c) as V_{ig} and V_{bg} , respectively. The large gold pad at the upper-left corner of the figure is the ionic liquid side-gate electrode. The ratio of the side-gate area to the BLG area should be maximized in order to achieve optimum ionic liquid gating effect. In our device shown in Fig. 1(c,d), the ratio of the side-gate area to the effective BLG area within the opening window is more than 10^3 , and thus the applied V_{ig} can be fully imposed upon the EDL at the surface of BLG.

The freezing point of DEME-TFSI ionic liquid is at about $T = 180$ K, and precaution is needed while cooling down the device near the freezing point to avoid possible damage due to the volume change of ionic liquid below the freezing point. In addition, since the ions are frozen below 180 K, one would need to warm up the ionic liquid to temperature higher than the freezing point to change the doping level. Figure 2(a) plots the sheet resistance R_{\square} versus V_{bg} at three different ionic liquid side gate voltages V_{ig} values. The device was first warmed up to $T = 220$ K followed by setting a new V_{ig} value. Due to the slow motion of ions, we typically waited for at least 30 minutes till the ions reached steady state, and then it was slowly cooled down to $T = 20$ K for resistivity and TEP measurements. The progressive shift of the charge neutral point (CNP) to higher V_{bg} values as V_{ig} goes from $+1.5$ V to -1.5 V is an indication of effective charge doping in the BLG. Both the peak location and also R_{\square} at CNP shows nearly linear dependence on the V_{ig} as demonstrated in Fig. 2(b). The linear fit to the data points give

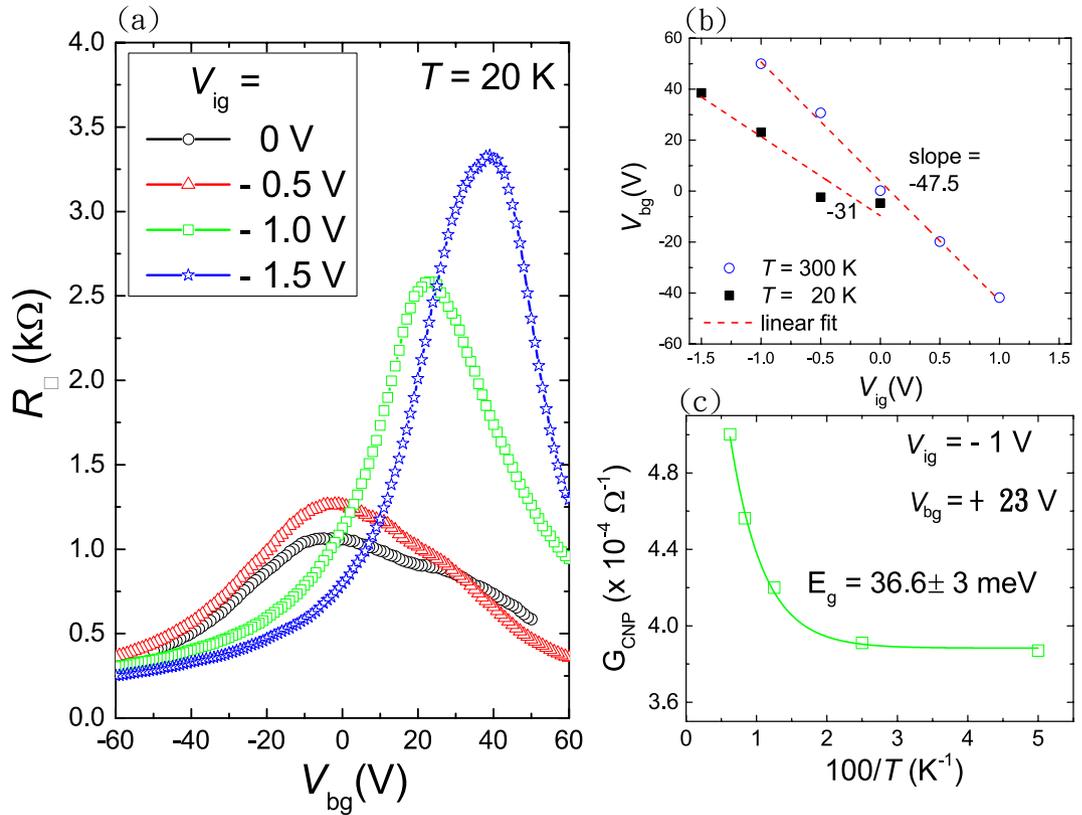


Figure 2. (a) Sheet resistance R_{\square} versus bottom gate V_{bg} at $T = 20$ K under three different side gate voltages V_{ig} . (b) The resistance peak location in (a) as a function of side gate voltage V_{ig} at $T = 300$ K and 20 K. (c) Conductance at CNP G_{CNP} versus $100/T$. A band gap of $E_g = 36.6 \pm 3$ meV is obtained for $V_{ig} = -1$ V. Solid line is the fitting curve based on the thermal-activation formula.

$dV_{bg}/dV_{ig} = -47.5$ at $T = 300$ K, which infers that the ionic liquid gating effectively introduces 47.5 times more charge carriers than regular bottom gate V_{bg} via SiO_2 (300 nm)/Si if the same voltage bias is used. At $T = 20$ K, dV_{bg}/dV_{ig} reduces to -31 , which suggests that nearly 30% of induced charges at the surface become immobile at lower T . We also notice that the full-width half maximum (FWHM) of the $R_{\square} - V_{bg}$ curve is getting smaller as the V_{ig} goes more negative in magnitude. This may imply a nontrivial charge doping effect via ionic liquid gating, where the charge puddle effects can be an important issue. From the T -dependence of conductance at CNP G_{CNP} (open square in Fig. 2(c)), a band gap of E_g about 36.6 ± 3 meV at $V_{ig} = -1$ V and $V_{bg} = +23$ V is obtained by fitting to the thermal activation formula $G_{CNP} = A \exp[-E_g/k_B T] + C$, where A and C are T independent constants. The resulting fitting curve is shown as solid line in Fig. 2(c).

In order to further look into the issue about FWHM in the $R_{\square} - V_{bg}$ curve, we carried out a series of experiments by cooling down the device at different V_{bg} values after setting a new V_{ig} value. Surprisingly, we found that the $R_{\square} - V_{bg}$ curve at low T is strongly dependent on the initial state before cooling. When the Fermi level was kept at CNP during the device cool-down, the $R_{\square} - V_{bg}$ curve shows much larger R_{\square} at CNP and smaller FWHM compared to the one with Fermi level away from the CNP during the cool-down. This is demonstrated in Fig. 3, where all the three curves are measured at $V_{ig} = -1.5$ V. The black curve is the data at $T = 220$ K above the DEME-TFSI freezing point, which indicates that Fermi level can be shifted to the CNP by applying a $V_{bg} = +43$ V. The blue and red curves are the $R_{\square} - V_{bg}$ data at $T = 20$ K with application of $V_{bg} = +43$ V and 0 V during the cool-down, respectively. Remarkably, the blue curve (open circle symbol) shows a FWHM of about 30 V and a large peak at $V_{bg} = +38$ V while the red curve (open square symbol) only shows a gradual increase in R_{\square} with increasing V_{bg} up to $+70$ V without attaining a local maximum. This finding suggests that the charge puddles condition in BLG is strongly influenced by the ionic liquid, which requires careful control over the cooling procedure. One likely explanation is due to the non-uniform charge distribution in the EDL on the surface of BLG, which is likely due to the defected condition of the interface between the BLG and ionic liquid. By setting the Fermi level at CNP during the cool-down, it minimizes the charge screening effect with the mobile ions and thus introduces smaller charge puddles at low T resulting in a better device performance.

By using the same cool-down sequence of setting Fermi level at CNP, we also performed TEP measurement on a BLG device with ionic liquid gating. The local heater current was modulated at a low frequency ω , and the TEP signals were taken by a lock-in amplifier at second harmonic 2ω . The temperature gradient across the BLG was determined by two local thermometers made of thin gold nanowires. The resulting TEP S versus V_{bg} at $V_{ig} = -1$ V and 4 different temperatures is shown in Fig. 4(a), where the definitions for the separation of the local extremes ΔV_{bg} and TEP S_m are given. Both the $S - V_{bg}$ and the $R_{\square} - V_{bg}$ curves at lower temperatures show nearly no

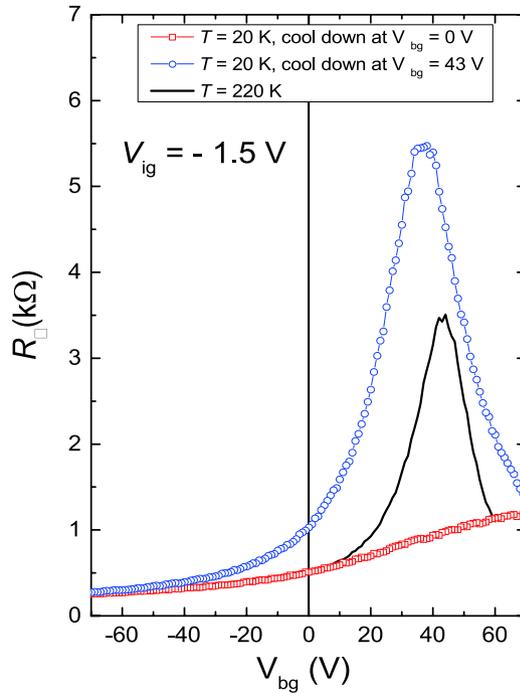


Figure 3. Sheet resistance R_{\square} versus bottom gate V_{bg} at $T = 20$ K and $V_{ig} = -1.5$ V. Blue and red data points are the data taken at $T = 20$ K after cooling down the device at $V_{bg} = 0$ V and $+43$ V, respectively. The solid black line is the $R_{\square} - V_{bg}$ curve at $T = 220$ K, which is above the freezing point of the ionic liquid DEME-TFSI.

hysteresis when sweeping V_{bg} . When Fermi level is tuned to the CNP, the TEP is nearly zero. The TEP reaches local extreme values as the Fermi level shifts away from the CNP. We note that all the $S - V_{bg}$ curves shows only single sudden drop in $S = 0$ at $\sim V_{bg} = +23$ V suggesting that most of the thermoelectric signals come from the ionic liquid gating effective region. The contribution from non-effective regions outside of the opening window shown in Fig. 1(d) with CNP at $V_{bg} \cong 0$ V is thus negligible. On the other hand, the ΔV_{bg} is about 40 V consistent with the FWHM of the $R_{\square} - V_{bg}$ curve shown in Fig. 3. The maximum TEP $S_m/2$ grows with increasing temperature as shown in Fig. 4(b), where $S_m/2 - T$ curves at $V_{ig} = 0, -0.5$, and -1 V are plotted. A clear enhancement of the TEP is observed when applying ionic liquid side gate V_{ig} . At $T = 120$ K, the TEP $S_m/2$ increases by about 40% from $30 \mu\text{V/K}$ at $V_{ig} = 0$ to $42 \mu\text{V/K}$ at $V_{ig} = -1$ V.

The enhancement of TEP in BLG device using ionic liquid does not give a larger TEP compared to the dual-gated BLG device using hBN as top gate dielectric⁹. One important factor is the electron-hole puddles effect, which turns out to be more significant in ionic liquid gating device as pointed out earlier. This can be seen more clearly from the variation of ΔV_{bg} at different V_{ig} values shown in Fig. 4(c). ΔV_{bg} equals about 45 V at $V_{ig} = 0$ and systematically drops to 20 V at $V_{ig} = -1$ V, which remains bigger than other dual-gated BLG devices typically showing $\Delta V_{bg} \leq 10$ V^{9,16}. It is known that the electrons and holes both contribute to the total conductivity, but they would cause a compensation effect to the TEP signal^{17–20}. By taking a simplified two-band model with both electron and hole pockets at the Fermi level⁹, the total TEP can be expressed as $S = S'(|n_h| - |n_e|)/(|n_h| + |n_e|)$, where S' and $n_{e(h)}$ are the TEP for each band and the electron (hole) density, respectively. When the Fermi level is at the CNP, the total S equals zero. TEP S starts to increase rapidly as the Fermi level shifts away from the CNP, and it reaches a local extreme when one of the pockets was fully depleted. Therefore, the density n_m at the local extremes in $S - V_{bg}$ curve is a good indication of the amount of charge puddles near the CNP in a BLG device. n_m can be estimated using the formula of $n = C_g/e V_{bg}$, where C_g is the gate capacitance and equals $\cong 140 \text{ aF}/\mu\text{m}^2$ determined from the Hall effect measurements. As shown in Fig. 4(c), n_m equals $1.0 \times 10^{12} \text{ cm}^{-2}$ at $V_{ig} = -1$ V, which is more than two-fold larger than that in a dual-gated BLG device using hBN as top gate dielectric. As suggested by the theoretical calculations^{9,21}, the magnitude of TEP enhancement is largely dependent on n_m , where significant TEP enhancement can be achieved in devices with much lower n_m values. The charge puddles effect near CNP appears to be less significant at higher V_{ig} values, which is likely associated with the improved charge distribution in the EDL on the surface of BLG under finite electric bias.

In summary, we have demonstrated the feasibility of using ionic liquid as a top gate to induce a band gap in a BLG. We found that the cooling procedure of the device can strongly affect the amount of charge puddles in a BLG possibly associated with the non-uniform charge distribution in the EDL of ionic liquid. A lowest level of charge puddles can be achieved by cooling down the device while keeping the Fermi level at the CNP. In this way, we successfully demonstrate the enhancement of TEP by nearly 40% with a merely -1 V application to the ionic liquid side gate. For further enhancement in TEP solution to the reduction of charge puddles near CNP and a cleaner graphene surface condition with proper device cooling procedure are both needed and require more experimental efforts.

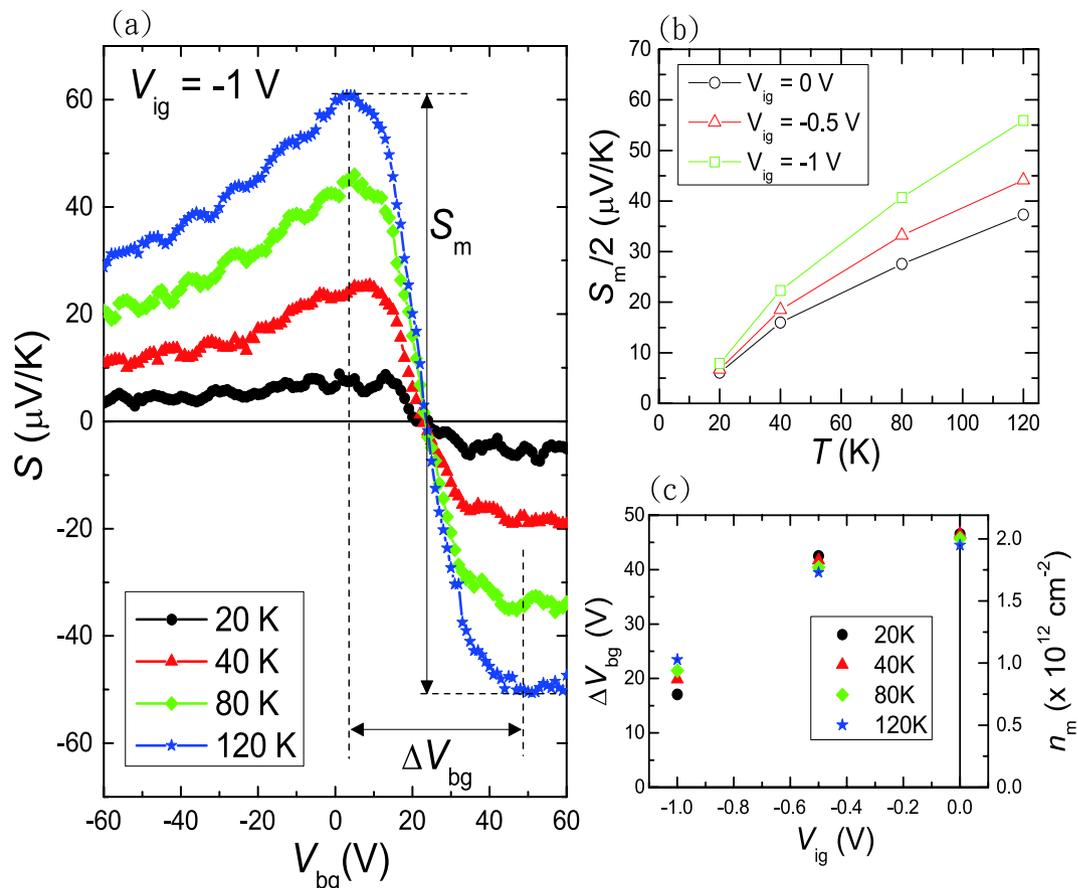


Figure 4. (a) TEP S versus V_{bg} at $V_{ig} = -1$ V at 4 different temperatures. (b) The temperature dependence of $S_m/2$ at three different side gate voltages of $V_{ig} = 0, -0.5$ and -1 V. (c) ΔV_{bg} and the corresponding n_m as a function of V_{ig} . A progressive decrease in ΔV_{bg} with increasing magnitude of V_{ig} is shown.

References

1. Qiao, Z., Tse, W. K., Jiang, H., Yao, Y. & Niu, Q. Two-dimensional topological insulator state and topological phase transition in bilayer graphene. *Phys. Rev. Letts* **107**, 256801 (2011).
2. Maher, P. *et al.* Evidence for a spin phase transition at charge neutrality in bilayer graphene. *Nature Phys.* **9**, 154 (2013).
3. Taychatanapat, T., Watanabe, K., Taniguchi, T. & Jarillo-Herrero, P. Electrically tunable transverse magnetic focusing in graphene. *Nature Phys.* **9**, 225 (2013).
4. Ling, X. *et al.* Can graphene be used as a substrate for Raman enhancement? *Nano Letts.* **10**, 553 (2009).
5. Grigorenko, A. N., Polini, M. & Novoselov, K. S. Graphene plasmonics. *Nature Photon.* **6**, 749 (2012).
6. Garcia de Abajo, F. J. Graphene nanophotonics. *Science* **339**, 917 (2013).
7. Ohta, T., Bostwick, A., Seyller, T., Horn, K. & Rotenberg, E. Controlling the electronic structure of bilayer graphene. *Science* **313**, 951 (2006).
8. Zhang, Y. *et al.* Direct observation of a widely tunable bandgap in bilayer graphene. *Nature* **459**, 820 (2009).
9. Wang, C. R. *et al.* Enhanced thermoelectric power in dual-gated bilayer graphene. *Phys. Rev. Letts* **107**, 186602 (2011).
10. Cutler, M. & Mott, N. F. Observation of Anderson localization in an electron gas. *Phys. Rev.* **181**, 1336 (1969).
11. Fallahzad, B., Kim, S., Colombo, L. & Tutuc, E. Dielectric thickness dependence of carrier mobility in graphene with HfO_2 top dielectric. *Appl. Phys. Letts* **97**, 123105 (2010).
12. Ye, J. T. *et al.* Liquid-gated interface superconductivity on an atomically flat film. *Nature Mater.* **9**, 125 (2010).
13. Bollinger, A. T. *et al.* Superconductor-insulator transition in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ at the pair quantum resistance. *Nature* **472**, 458 (2011).
14. Ye, J. T. *et al.* Accessing the transport properties of graphene and its multilayers at high carrier density. *Proc. Natl. Acad. Sci. USA* **108**, 13002 (2011).
15. Yamashiro, Y., Ohno, Y., Maehashi, K., Inoue, K. & Matsumoto, K. Electric-field-induced band gap of bilayer graphene in ionic liquid. *J. Vac. Sci. Technol. B* **30**, 03D111 (2012).
16. Taychatanapat, T. & Jarillo-Herrero, P. Electronic transport in dual-gated bilayer graphene at large displacement fields. *Phys. Rev. Letts* **105**, 166601 (2010).
17. Adam, S., Hwang, E. H., Galitski, V. M. & Das Sarma, S. Self-consistent theory for graphene transport. *Proc. Natl. Acad. Sci. USA* **104**, 18392 (2007).
18. Xiao, S., Chen, J. H., Adam, S., Williams, E. D. & Fuhrer, M. S. Charged impurity scattering in bilayer graphene. *Phys. Rev. B* **82**, 041406(R) (2010).
19. Das Sarma, S., Hwang, E. H. & Rossi, E. Theory of carrier transport in bilayer graphene. *Phys. Rev. B* **81**, 161407(R) (2010).
20. Hwang, E. H., Rossi, E. & Das Sarma, S. Theory of thermopower in two-dimensional graphene. *Phys. Rev. B* **80**, 235415 (2009).
21. Hao, L. & Lee, T. K. Thermopower of gapped bilayer graphene. *Phys. Rev. B* **81**, 165445 (2010).

Acknowledgements

The authors acknowledge the funding support from Ministry of Science and Technology (MOST) in Taiwan (NSC102-2112-M-001-022-MY3) and technical support from Core Facility for Nanoscience and Nanotechnology at Academia Sinica in Taiwan.

Author Contributions

W.-L.L. and H.Y. designed the experiment. Y.-Y.C. and C.-R.W. fabricated the BLG device and performed the low temperature transport measurements. W.-L.L. wrote the manuscript.

Additional Information

Competing financial interests: The authors declare no competing financial interests.

How to cite this article: Chien, Y.-Y. *et al.* Thermoelectric Power in Bilayer Graphene Device with Ionic Liquid Gating. *Sci. Rep.* **6**, 20402; doi: 10.1038/srep20402 (2016).



This work is licensed under a Creative Commons Attribution 4.0 International License. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons license, users will need to obtain permission from the license holder to reproduce the material. To view a copy of this license, visit <http://creativecommons.org/licenses/by/4.0/>